

What are so-called elementary particles?

N.B. All used equations use [N.U.](#), i.e. $h/(2\pi)$ ([Dirac's constant](#)) = $c = 1$

Elementary particles are a physical idea used by physicists to specify all non-reducible mathematical representations of all always-valid symmetry groups. But almost all physicists don't look at these particles this way. Instead, [elementary particles](#) are seen as not build from more other particles investigated particles.

The [S\(pecial\)R\(elativity\) Poincaré group](#) gives the most general continuous symmetry group after inclusion of gravitational effects like curvature and all appearing discrete symmetry groups like charge inversion or helicity flip in case of uncharged [fermions](#). All basic constituents of our universe are given completely with the non-reducible mathematical representations of all always-valid symmetry groups. I.e. particles with other characteristics then all characteristics resulting from this non-reducible complete symmetry analysis, must be fictitious, like for example the assumed very massive scalar (spinless) [Higgs boson](#)!

To start-off with, a relativistic 4D-spacetime description will be used. In [4] a simple proof is given to show that any non-3D-space description of any possible reality is bogus, i.e. cannot be used to describe [CAP](#) extended [fermions](#).

The only possible exact description of [elementary particles](#) must use a linear algebra, i.e. use a [mathematical](#) description in [linear space](#), such that differentials and integrals analyzed with summed infinitesimal steps are possible.

The used description must be relativistic. Each observer has only one possible maximum speed with always the same constant value, called the light-speed, independent of the motion of the observer with respect to any possible inertial frame. In [SR](#) one only analyzes relativistic space-time in so-called [Minkowski-space](#) and first order time derivatives of 4-vectors, in other words accelerations are not actually analyzed. [QM](#), i.e. our microscopic world is only analyzed [SR](#). At first sight this linear analysis seems correct, however it also appears that this mathematical analysis does not conform to the [CAP](#) so can **NOT** be correct! However, it appears to be easy to re-write [QM](#) such that it complies to the [CAP](#).

To re-write the [S\(tandard\) M\(odul\) of SR Q\(uantum\)F\(ield\)T\(theories\)](#) to comply to the [CAP](#) we should start with a short analysis of [G\(eneral\)R\(elativity\)](#).

In [GR](#) one first shows that inert mass is equivalent to heavy mass, i.e. that acceleration of observers is identical to experiencing weight resulting from earth gravity. In the end the description is adapted, such that the description is valid for any possible descriptive frame, instead of just the inertial frames used in the theory of [SR](#). In the used explanation [1] is used throughout. The mathematical solution of the equations of motion of [GR](#) shows curvature of space-time. The equations of motion must be invariant after coordinate transformations. This demands all derivatives in the equations of motion to be so-called co-variant derivatives, given by “:”.

For example, the co-variant derivative of a 4-vector A_μ is given by:

$$A_{\mu;\nu} \equiv \frac{\partial A_\mu}{\partial x^\nu} \rightarrow A_{\mu;\nu} \equiv A_{\mu,\nu} - \Gamma^\sigma_{\mu\nu} A_\sigma, \text{ with Christoffel symbol } \Gamma^\sigma_{\mu\nu} = g^{\sigma\alpha} \Gamma_{\alpha\mu\nu} = g^{\sigma\alpha} \left\{ \frac{1}{2} (g_{\alpha\mu,\nu} - g_{\mu\nu,\alpha} + g_{\nu\alpha,\mu}) \right\} \quad (1)$$

Ordinary derivatives always commute, co-variant derivatives do in general not commute!

Besides (1) all matrices, must be so-called tensors, i.e. transform according to:

$$T^{\alpha'\beta'}_{\gamma'} = x^{\alpha'}_{,\lambda} x^{\beta'}_{,\mu} x^{\nu}_{,\gamma'} T^{\lambda\mu}_\nu \quad (2)$$

According to [Einstein's CAP](#), every description of physics must include the gravitational action, because all other actions **always** depend on the gravitational action. The fact that co-variant derivatives don't commute results in a description, outside always-possible singularities (terminating points of [black holes](#) and the so-called [Big Bang](#)), with degrees of freedom **doubled** [3]. The gravitational-action alone is given with the symmetrical [Ricci tensor](#) and the curvature scalar R multiplied by the [fundamental tensor \(metric\)](#) $g_{\mu\nu}$:

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 0 \quad (3)$$

The [Ricci tensor](#) follows from the [Riemann-Christoffel](#) tensor, i.e. the so-called curvature-tensor with 4 indices and given by:

$$R_{\mu\nu\alpha\beta} = \Gamma_{\mu\nu\beta,\alpha} - \Gamma_{\mu\nu\alpha,\beta} + \Gamma_{\nu\beta}^{\gamma}\Gamma_{\mu\gamma\alpha} - \Gamma_{\nu\alpha}^{\gamma}\Gamma_{\mu\gamma\beta} \quad (4)$$

This tensor complies with the [Bianchi](#) symmetry relations and as a result of that only has 20 independent degrees of freedom. Exactly the degrees of freedom of the [Ricci tensor doubled](#), which follows from the curvature tensor contracting two indices, such that the result isn't zero. Outside singularities curvature can be described in a so-called linear Riemann-space with degrees of freedom [doubled](#), see [1] equation (6.1) and [3]. All Riemann degrees of freedom, which can't be interpreted, are removed from all resulting expressions in [GR](#) with the so-called [Christoffel symbol](#). This is why the curvature tensor given in (4) is given completely with [Christoffel symbols](#).

Symmetry groups are transformation groups, which result into invariant equations of motion. Such transformations are either continuous or discrete.

The complete continuous symmetry group of [GR](#) is the [extended Poincaré-group](#), i.e. the 6D-[Lorentz-group](#) together with all 4D-spacetime translations, and with spin2 [gravitational-field](#) induced symmetries, i.e. [curvature of space-time](#).

[Discrete symmetry](#) transformations can be used from the [SM](#), because the spin2 symmetrical [gravitational-field](#) only describes attractive forces, so there are no change-of-sign symmetries. For example charge-inversion or parity transformation (parity-inversion).

All symmetry transformations of the [CAP-compliant extended Poincaré-group](#) can be given with the sum of an anti-symmetrical transformation tensor $A_{\mu\nu}$ and a symmetrical transformation tensor $S_{\mu\nu}$:

$$T_{\mu\nu} = A_{\mu\nu} + S_{\mu\nu}, \text{ with } A_{\mu\nu} = -A_{\nu\mu} \wedge S_{\mu\nu} = S_{\nu\mu} \quad (5)$$

This representation in principle has 16 degrees of freedom, but all [extended Poincaré-group](#) transformations are either symmetrical, or anti-symmetrical, such that any arbitrary [extended Poincaré](#) transformation can always be represented by (5). This analysis gradually shows how a mathematical description of all elementary particles comes to life.

All possible particles with all their characteristics follow uniquely from a non-reducible description of the most general symmetry group analysis. In our 4D-spacetime universe the most general symmetry groups are the continuous extended Poincaré-group (5) and all discrete symmetries, which can be analyzed independently.

The [gravitational-field](#) given by (3), obviously is a symmetrical action. The gravitational field has mass as primary source. All other actions, either bose- or fermi-like, all depend on the gravitational action and give contributions in equation (3), i.e. the zero must be replaced by all actions which interact with the gravitational field. All contributions of these actions are symmetrical too.

[Albert Einstein](#) derived the following equations of motion after including all possible actions. This equation follows from an [Euler-Lagrange](#) description in which the constant multiplied with infinitesimal delta $\delta g_{\mu\nu}$ must be zero:

$$R_{\mu\nu} - (\frac{1}{2}R + \lambda)g_{\mu\nu} - 8\pi(\rho v_{\mu}v_{\nu} + E_{\mu\nu}) = 0 \Leftrightarrow R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = 8\pi(\rho v_{\mu}v_{\nu} + E_{\mu\nu}) + \lambda g_{\mu\nu} \quad (6)$$

Also see [1], equation (29.7) with added the freedom of the [cosmological constant](#).

With ρ the mass density and $v_{\mu} = \frac{\partial x_{\mu}}{\partial t}$ the speed 4-vector of mass density at point x_{μ} .

Constant λ is the so-called [Cosmological constant](#), which describes a spin 0 action, i.e. a spinless action, which isn't possible because it does **NOT** comply to the [CAP](#), also see [3]!

$E_{\mu\nu}$ is the stress-energy tensor of the [EM-field](#) and is related to the anti-symmetrical tensor $F_{\mu\nu}$ build from the space-like electrical field \mathbf{E} (fat notation represents space-like 3D-vectors) and magnetic field \mathbf{B} :

$$F_{0i} = (E_x, E_y, E_z) \wedge F_{23} = B_x \wedge F_{31} = B_y \wedge F_{12} = B_z \quad (7)$$

Also see [2], tensor (5.1).

Tensor $F_{\mu\nu}$ is anti-symmetrical, so the [Maxwell equations](#) are easily given invariant by:

$$F_{\alpha\beta;\gamma} + F_{\beta\gamma;\alpha} + F_{\gamma\alpha;\beta} = F_{\alpha\beta;\gamma} + F_{\beta\gamma;\alpha} + F_{\gamma\alpha;\beta} = 0 \quad (8)$$

The stress-energy tensor of the [ElectroMagnetic-field](#) is given by the anti-symmetrical tensor $F_{\mu\nu}$, i.e. the [EM-fields](#) according to:

$$E_{\mu\nu} = -\frac{1}{4\pi} F_{\mu}^{\alpha} F_{\alpha\nu} + \frac{1}{16\pi} g_{\mu\nu} F^{\alpha\beta} F_{\alpha\beta} \quad (9)$$

The [EM-fields](#) follow uniquely after applying [Lorenz-gauge](#) symmetry from the [vector-potential](#) $A_{\mu} = (V, \mathbf{A})$ of the [EM-field](#):

$$F_{\mu\nu} = A_{\mu;\nu} - A_{\nu;\mu} = A_{\mu;\nu} - A_{\nu;\mu} \quad (10)$$

The δA_{μ} term gives: $F_{\mu;\nu} = 4\pi J_{\mu}$ (11)

Here the current density is given by $J_{\mu} = \sigma v_{\mu} \sqrt{}$, with $\sqrt{} \equiv \sqrt{(-\det(g_{\mu\nu}))}$ (12)

With σ the charge density in Einstein's mathematical continuously analyzed description. Equations (10) and (11) represent the [Maxwell equations GR](#), which give the [EM-field](#) completely after imposing the [SR Lorenz-gauge](#) symmetry.

All in experiments observed charge always has mass. This experimental fact is valid for all charged particles, whether [bosons](#) or [fermions](#). This is because the spin1 [EM-field](#) allows interaction with charge in all directions, just like the spin2 [gravitational field](#) interacts with mass in all directions.

The mass-sources of the [gravitational-field](#) can be given with mass current density, i.e. momentum-density, just like (12):

$$p_{\mu} = \rho v_{\mu} \sqrt{} \quad (13)$$

The last delta-contribution δx_{μ} shows that charged mass experiences the so-called [Lorentz-force](#) and as a result follows a path different from the [GR geodesic](#) path:

$$\rho v^{\nu}{}_{;\mu} v_{\nu} + F_{\mu}{}^{\nu} J_{\nu} = 0 \quad (14)$$

On one hand Einstein analyzed everything with a model with continuous space-time and continuous divided mass- and charge-densities, on the other hand he already drew the conclusion in 1905 that all [photons](#) are elementary particles with energy proportional to the frequency of the carried [EM-field](#). So why he analyzed everything in his [GR](#) with continuous varying densities remains a mystery to me!?!

[Photons](#) are the elementary particles, which together represent the [EM-field](#). In other words the [photon](#) is nothing but the mathematical non-reducible representation of the [EM-field](#), i.e. the 4D-spacetime [vector-potential](#) $A_{\mu} = (V, \mathbf{A})$. This particle appears to be a spin1 particle, and is necessary to describe the anti-symmetrical symmetry group of the [extended Poincaré-group](#) non-reducible. In this way the so-called [extended Poincaré-group](#) is rewritten in compliance with both [SR](#) and Einstein's [CAP](#). This will be explained just before (26).

Up to now the relativistic analysis is described using continuous densities. In this description mass (13) and charge (12) are analyzed exactly the same.

However, there's a big difference between the anti-symmetrical [EM-field](#) and the symmetrical [gravitational-field](#):

A coordinates transformation $x \rightarrow x'$ can be described vector-like as: $d^4x' = J d^4x$ (15)

With J the [Jacobian](#), i.e. the determinant of $x'^{\mu}{}_{;\nu}$.

The [metric](#) transforms as: $g_{\alpha\beta} = x'^{\mu}{}_{;\alpha} x'^{\nu}{}_{;\beta} g_{\mu\nu}$ (16)

This equality is also valid for the determinants: $g = J^2 g'$, with $g = \det(g_{\mu\nu}) < 0$ (17)

The square-root of (17), with definition (12), gives for the change: $\sqrt{} = J \sqrt{}'$ (18)

All equations of motion must be invariant after arbitrary coordinate transformations.

A scalar $S = S'$, only has the following invariant space-time integral:

$$\int S \sqrt{d^4x} = \int S' \sqrt{J} d^4x = \int S' \sqrt{J'} d^4x' \quad (19)$$

This is why \sqrt{J} is called a scalar density. This is also why the square-root of the absolute value of the determinant of the metric is very important in [GR](#).

Tensors with one or more indices are analogously defined as a tensor-density after multiplication with \sqrt{J} . Integrated tensor densities are in principle **NOT** conserved quantities, because a tensor integrated over a larger space-time area doesn't transform linearly after an arbitrary coordinates transformation due to curvature. Only the self-contracted co-variant divergence of a 4-vector always results in the [GR Gauss theorem](#):

$$V^{\mu}_{;\mu} = V^{\mu}_{;\mu} - \Gamma^{\mu}_{\mu\nu} V^{\nu} = V^{\mu}_{;\mu} - \sqrt{-1} \sqrt{J}_{,\nu} V^{\nu} \Leftrightarrow V^{\mu}_{;\mu} \sqrt{J} = (V^{\mu} \sqrt{J})_{;\mu} \quad (20)$$

I.e. $V^{\mu}_{;\mu}$ is a scalar, just like $S = S'$ in (19).

The co-variant derivative of a tensor with 2 or more indices doesn't result in an integrated scalar and in general can't be analyzed with the [Gauss theorem](#) (20).

There is one exception, the **anti-symmetrical** tensor $F^{\mu\nu}$: $F^{\mu\nu}_{;\nu} = F^{\mu\nu}_{;\nu} + \Gamma^{\mu}_{\nu\alpha} F^{\alpha\nu} + \Gamma^{\nu}_{\nu\alpha} F^{\mu\alpha} = F^{\mu\nu}_{;\nu} + \sqrt{-1} \sqrt{J}_{,\alpha} F^{\mu\alpha}$ (21)

So, just as in the case of the scalar S in (19): $F^{\mu\nu}_{;\nu} \sqrt{J} = (F^{\mu\nu} \sqrt{J})_{;\nu}$ (22)

When $F^{\mu\nu}_{;\nu} = 0$, one yields a conservation law of a continuous vector density $F^{\mu 0} \sqrt{J}$. If no net-current passes a selected area covering the integrated space, then $\int F^{\mu 0} \sqrt{J} d^3x$ is a constant.

The 2 indices symmetrical tensor in (5) results in an additional contribution:

$$S_{\mu;\nu} = S_{\mu;\nu} + \sqrt{-1} \sqrt{J}_{,\alpha} S_{\mu}^{\alpha} - \Gamma_{\alpha\mu\nu} S^{\alpha\nu} \quad (23)$$

$S^{\alpha\nu}$ is symmetrical, so the last term can be written as: $\frac{1}{2}(\Gamma_{\alpha\nu\mu} + \Gamma_{\nu\alpha\mu}) = \frac{1}{2}g_{\alpha\nu,\mu}$ (24)

I.e. the symmetrical 2 indices tensor complies to: $S_{\mu;\nu} \sqrt{J} = (S_{\mu}^{\nu} \sqrt{J})_{;\nu} - \frac{1}{2}g_{\alpha\beta,\mu} S^{\alpha\beta} \sqrt{J}$ (25)

Also see [1], equation (21.4).

The spin 1 [EM-field](#) is described with an anti-symmetrical field tensor (10). The equations of motion and conservation laws are solved relatively easy in this case. In the case of the [gravitational-field](#) given with the symmetrical [metric \(fundamental tensor\)](#), via difference-terms of derivatives (like given in (1)), such that the complete expression is a co-variant derivative, has a nonlinear character due to curvature and as a result of that can't be analyzed with the [Gauss theorem](#). Also compare (22) with (25).

According to (5) the [CAP extended Poincaré-group](#) is given completely with the sum of a symmetrical and an anti-symmetrical transformation tensor, with respectively 10 and 6 degrees of freedom. When also taking into account the fact that the graviton is a spin2 particle this results in the 10 degrees of freedom of the symmetrical transformation tensor being **doubled** consistent with compliance to the [CAP](#).

From these non-reducible degrees of freedom all so-called [elementary particles](#) come to life together with all their characteristics in a mathematical description.

From [QM](#) it is well known that all elementary particles have [intrinsic angular momentum](#), the so-called [spin](#).

With simple analysis [3] I became convinced that all elementary particles must be described extended in the 2D-plane orthogonal to the observed direction of motion. Any not extended elementary particle can only be on the [SR wordline](#), so it can't have any wavelike character specified by a frequency and doesn't comply to Einstein's [CAP](#)! The [CAP](#) imposed extendedness in the 2D-plane orthogonal to the observed direction of motion ([wordline](#)) has in the description from the inertial frame with origin at the average position of the oscillating point constants of motion. The constant time-like component just is the total energy, which must be equal to the detected frequency of the oscillating point multiplied by [Planck's constant](#). The space-like 3D-vector just is the constant angular momentum, which of course just is the particle's [spin](#) given by the half integer spin s multiplied by Dirac's constant, i.e. [Planck's constant](#) divided by 2π . The average extendedness in polar coordinates described from the inertial frame with origin at the average position is proportional with the [spin](#) (actually the [helicity](#)) multiplied by the Golden Ratio: $2\langle\rho\rangle = \text{spin} \times \text{Golden Ratio} \times \text{Planck length}$ (26)

First of all this explains why all [elementary particles](#) have sizes of the order of the [Planck length](#). Second, **spinless elementary particles** do **NOT** possess energy proportional to a frequency so are simple human fiction.

And last but not least, the average extendedness is proportional to the [Golden Ratio](#) and this constant explains why this ratio occurs so often in everyday life.

Both the 10 and 6 degrees of freedom of the symmetrical and anti-symmetrical transformation sets must be represented by a half-integer [spin](#), which represents the source, and an integer [spin](#) that represents the resulting force field.

I.e. the complete continuous [CAP extended Poincaré-group](#) (5) must be represented as:

$$\text{spin}\frac{1}{2} \otimes \text{spin}2 \oplus \text{spin}\frac{1}{2} \otimes \text{spin}1 \quad (27)$$

N.B. a spinless particle with one degree of freedom (scalar) is not used, because an elementary spinless particle never leaves its average [SR worldline](#) and as a result of that doesn't comply with the [CAP](#) also see [\[3\]](#).

The symmetrical set describes the stable [spin](#) $\frac{1}{2}$ masses multiplied by the resulting [spin](#)2 [gravitational-field](#).

The anti-symmetrical set gives the stable [spin](#) $\frac{1}{2}$ charges multiplied by the resulting [spin](#)1 [EM-field](#).

The [EM-field](#) is described with the 4D-[vector-potential](#) $A_\mu=(V, \mathbf{A})$.

A representation with spin s has $(2s+1)$ degrees of freedom. The spin-representation of the [extended Poincaré-group](#) (26) shows that elementary particles with $\text{spin} > 2$ are not allowed in our 4D-spacetime universe. With [\[4\]](#) I have shown that [knots](#) are only possible in 3D-space, however [Grisha Perelman](#) has proven this in 2003 at [Stony Brook university](#) in New York together with Richard Hamilton. His work is given in [\[5\]](#), [\[6\]](#) and [\[7\]](#). The mathematical solution of the extendedness of all elementary particles in the 2D-plane orthogonal to the average direction of motion, described by the relativistic worldline is a solution of [SR D\(ifferential\)E\(quations\)](#). To give this solution completely one must impose [B\(oundary\)C\(onditions\)](#). The searched for solutions here have 2 possible independent BC, open BC and closed BC.

Open BC describe elementary particles which are able to interact with other particles in all directions.

Closed BC describe elementary particles which only interact in the direction of motion, i.e. head on which coincides with the traveled [SR wordline](#). But this [wordline](#) now can only be described linear at infinitesimal scale, so curvature of space and time apparently shows itself on two different grounds. And any valid [SR](#) description still has to comply with the [CAP](#), i.e. has all non-reducible mathematical representations (i.e. [elementary particles](#)) described as extended harmonic oscillating points in the 2D-plane orthogonal to the [wordline](#).

Open BC have a positive integer as degree of freedom extra. This must be the [quantum number](#) of the described particles family. The higher this number the more interaction with the [gravitational-field](#), so the higher the mass. This is why open BC describe [fermions](#) that apparently always have mass. The closed BC consequently describe [bosons](#), of which only one fundamental specie exists of all possible symmetry groups. Only [bosons](#) are possibly massless. The fundamental force particles, which represent the [EM-field](#) and the [gravitational-field](#) non-reducible, are the [spin](#)1 [photon](#) and the [spin](#)2 [graviton](#) respectively. Consequently these are the only two possible massless particles!

In a [SR](#) analysis a path of an harmonic oscillating point is traceable in the 2D-plane orthogonal to the average position given by the [wordline](#). Fermions, always have mass, i.e. can always be described with knots in the traveled harmonic oscillating path. I'm not saying it'll actually happen, only perhaps near singularities, however mathematical it's always possible! This fact made me draw the conclusion that [fermions](#) are only possible in 4D-spacetime, in which [Albert Einstein](#) developed his theories of relativity.

Without [fermions](#) also no [bosons](#), i.e. nothing at all!

This is why all possible [universes](#) must have 4D-spacetime!

According to representation (27) only the following elementary spins are possible: $s \in \{\frac{1}{2}, 1, 1\frac{1}{2}, 2\}$ (28)

The given [spin](#) $1\frac{1}{2}$ particles cannot exist on their own, because this representation is not given in representation (27) of the [CAP extended Poincaré-group](#). From this fact I concluded that all [quarks](#) aren't [spin](#) $\frac{1}{2}$ particles with so-called [isospin](#) $\frac{1}{2}$, but all [quarks](#) are just elementary [spin](#) $1\frac{1}{2}$ particles. All possible [hadrons](#) are 3 combined [quarks](#) with total [spin](#) $\frac{1}{2}$, of which 3 elementary particle families exist in our universe. All elementary quarks have both mass and charge.

The charge always is a fraction of the electron charge e : $e_q \in \{-\frac{2}{3}e, -\frac{1}{3}e, \frac{1}{3}e, \frac{2}{3}e\}$ (29)

The fractions are such that all [hadron](#) charges are $\{-1, 0, 1\}$ multiplied by the electron-charge. This is because the anti-symmetrical [EM-field](#) complies to (22). As a direct result any possible universe has one unique charge-constant for all relative stable particles, just like every [universe](#) has its own specific light-speed, amount of fermion families, etcetera. For mass this characteristic isn't possible because the symmetrical spin2 action (25) doesn't comply to the [Gauss theorem](#).

The [elementary](#) spin $\frac{1}{2}$ particles are the so-called [leptons](#), of which also 3 families exist in our universe and are elementary particles with integer electron-charge $\pm e$, or chargeless and called [neutrino](#)'s. All [neutrino](#)'s always have rest-mass > 0 because they must be described with open BC. The [bosons](#) of (28) are not only the elementary spin1 [photon](#) and the elementary spin2 [graviton](#), but also all possible [gauge-bosons](#) available in a non-reducible mathematical analysis of any possible relativistic 4D-spacetime universe.

The [Lorenz-gauge](#) which specifies the [EM-field](#) completely with the [Maxwell equations](#) is a simple unitary one-dimensional U(1) [gauge-symmetry](#). The [EM-field](#) is described by the elementary U(1)-gauge massless spin1 [photon](#). In the only possible 4D-spacetime the maximum possible [gauge-symmetry](#) is just the gauge symmetry group used in the [SM](#): $U(1) \times SU(2) \times SU(3)$ (30)

The 'S' before the U(nitary) stands for the Special characteristic: $\det(SU(a)) = +1$ (31)
 With the Special property (31) all [gauge-symmetries](#), except the U(1) [gauge-symmetry](#), have their degrees of freedom correctly halved.

In a 4D-spacetime higher dimensional [gauge-symmetries](#) aren't possible:

$$\{(S)U(a) \mid 0 < a \in \text{integer} \wedge a < 4\} \quad (32)$$

The other gauge symmetries also imply particles. One should always remember that only the [anti-symmetrical](#) actions allow [gauge-symmetry](#).

The combined U(1) \times SU(2) [gauge-symmetry](#) group describes mixed by the so-called [Weinberg angle](#) the spin1 [photon](#) and the massive (but also uncharged) elementary [Z-gauge boson](#), which together with the charged (so massive) [W[±]-gauge bosons](#) represent the [weak nuclear forces](#) described by spin1 elementary particles. The SU(3) gauge symmetry group describes all spin $\frac{1}{2}$ [quarks](#), which combine in sets of 3 stable [quarks](#) kept together by [gluons](#) into stable spin $\frac{1}{2}$ so-called [baryons](#). Again 3 different particle-families exist in our universe with in the first family the +e charged [proton](#) and the uncharged, massive, [neutron](#). [Quarks](#) can also combine into [bosons](#), and are called [gluons](#) and [mesons](#). [Gluons](#) are combined [quark](#) and [anti-quark](#) pairs with a different (anti-) color. This characteristic allows [gluons](#) to keep [quarks](#) of a [baryon](#) together through so-called [color-exchange](#). This describes the so-called [strong nuclear force](#). [Mesons](#) are the remaining combined quark bosons, which are also possible besides the color-carrying [gluons](#).
 In principle it's possible that fermions build from more than 3 odd numbers of [quarks](#) exist, however in general such baryons are very unstable. In principle it's also possible that bosons build from more than 2 even numbers of [quarks](#) exist, however in general such mesons are again very unstable.

So, all possible [elementary particles](#) with all their characteristics follow completely from a non-reducible complete symmetries analysis of the only possible 4D-spacetime.

The complete continuous [CAP extended Poincaré-group](#) (27) has $2 \times 10 + 1 \times 6 = 26$ degrees of freedom instead of the only 10 degrees of freedom of the [SR](#) Poincaré symmetry group. Besides this complete continuous symmetry group the complete gauge symmetry group (30) must also be taken into account to end up with all possible elementary particles in any possible 4D-spacetime universe:

All possible elementary particles in our 3-particle families universe:

Fermions: 3 different families	Bosons: The elementary spin1 and spin2 bosons, see (26), and the U(1) x SU(2) x SU(3) gauge-bosons:
leptons: electron, muon and tauon + anti-particles	graviton, a spin2 elementary massless boson
leptons: massive but chargeless neutrino's	photon, a spin1 elementary massless boson
quarks 1st family: up-quark and down-quark	weak-nuclear forces: spin1 elementary massive gauge-bosons W [±] , Z
quarks 2nd family: charm-quark and strange-quark	strong-nuclear forces: spin1 colored quark+anti-quark gluons
quarks 3rd family: top-quark and bottom-quark	mesons: all non-gluon bose-quark combinations

All fermions have so-called [anti](#)-particles with changed charge sign of charged particles and opposite helicity in case of chargeless particles. All leptons are spin $\frac{1}{2}$ particles and all non-separable quarks are spin $\frac{1}{2}$ particles.

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